RESEARCH ARTICLE

# Kink and anti-kink wave solutions for the generalized KdV equation with Fisher-type nonlinearity

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## ABSTRACT

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# 1. Introduction

This study focuses on the travelling wave solutions of a newly introduced dispersion-convection-reaction model

$$u_t + \varepsilon u^n u_x + \mu u_{xxx} = ru(1 - u^n), \tag{1}$$

where n > 0,  $\varepsilon$  is a parameter for the purely nonlinear convection term,  $\mu$  is a parameter for the linear dispersion term and r is a parameter for the purely nonlinear reaction term. One can easily obtain the talented equations, such as the generalized KdV (gKdV) equation [1-9] and the dispersive-Fisher equation [10] by taking r = 0 and  $\varepsilon = 0$  in Eq. (1), respectively. There have been many cooperative combinations of dispersion with the different terms, such as dissipation, convection, diffusion and reaction [11-15]. Here, by combining the gKdV equation [2-8] with the Fisher-type (or KPP-type) nonlinearity [16-23], we call Eq. (1) as the gKdV-Fisher equation.

Recently, Galaktionov has focused on the higher-order versions of the KPP (or Fisher) type problem in the parabolic, dispersive and hyperbolic equations, see [21-23] for fruitful discussions. Fortunately, the third-order dispersive partial differential equation including the Fisher-type nonlinearity

$$u_{t} = u_{xxx} + u(1 - u), \tag{2}$$

with two travelling wave solutions

This paper proposes a new dispersion-convection-reaction model, which is called the gKdV-Fisher equation, to obtain the travelling wave solutions by using the Riccati equation method. The proposed equation is a third-order dispersive partial differential equation combining the purely nonlinear convective term with the purely nonlinear reactive term. The obtained global and blow-up solutions, which might be used in the further numerical and analytical analyses of such models, are illustrated with suitable parameters.

$$u_{1}(x,t) = \frac{1}{2} - \frac{9}{4} \tanh\left(\frac{19}{60}t - \frac{1}{60}9900^{\frac{1}{3}}(x+\xi_{0})\right) + \frac{11}{4} \tanh^{3}\left(\frac{19}{60}t - \frac{1}{60}9900^{\frac{1}{3}}(x+\xi_{0})\right),$$
$$u_{2}(x,t) = \frac{1}{2} + \frac{3}{4} \tanh\left(\frac{19}{60}t + \frac{1}{60}30^{\frac{2}{3}}(x+\xi_{0})\right) - \frac{1}{4} \tanh^{3}\left(\frac{19}{60}t + \frac{1}{60}30^{\frac{2}{3}}(x+\xi_{0})\right),$$

has been proposed in [10]. More recently, the exact travelling waves of the KdV-Burgers-Fisher equation

$$u_t + \varepsilon u u_x - v u_{xx} + \mu u_{xxx} = r u (1 - u), \qquad (3)$$

have been investigated in [24].

We would also like to remind the neighbouring nonlinear parabolic equation, which is a diffusionconvection-reaction model and called the generalized Burgers-Fisher equation [25,26],

$$u_t + \varepsilon u^n u_x - v u_{xx} = r u \left( 1 - u^n \right), \tag{4}$$

with the exact solution

$$u(x,t) = \left(\frac{\tanh(\xi)+1}{2}\right)^{\frac{1}{n}},$$

where

$$\xi = -\frac{n\varepsilon}{2\nu(n+1)}x + \frac{n\left(\varepsilon^2 + \nu r\left(n+1\right)^2\right)}{2\nu(n+1)^2}t + \xi_0.$$



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In the next section, we use the Riccati equation method [27-36] to reveal the travelling wave solutions of the gKdV-Fisher equation, which are the cooperative results of the proposed combined model. Here, it is reasonable to expect kink and antikink wave solutions because of the reaction term in the proposed equation.

# 2. Method and application

Let us first take the wave variable  $\xi = kx + wt + \xi_0$  and  $u(x,t) = u(\xi)$  in Eq. (1) to obtain the reduced nonlinear ODE as

$$wu' + \varepsilon ku^{n}u' + \mu k^{3}u''' - ru(1 - u^{n}) = 0.$$
 (5)

The solution of Eq. (5) is assumed to be expressed as

$$u = \sum_{i=0}^{M} a_i z^i , \qquad (6)$$

where the parameters,  $a_0, ..., a_M$  and M, can be determined later and  $z = z(\xi)$  is the solution of the following classical Riccati equation [27-35]:

$$z' = 1 - z^2, \tag{7}$$

which has the forms

$$z = \tanh(\xi) \text{ and } z = \coth(\xi).$$
 (8)

Here, using the advantage of the Riccati equation, higher-order derivatives of Eq. (7) can be obtained as

$$z'' = -2z + 2z^3, (9)$$

$$z''' = -2 + 8z^2 - 6z^4 \,. \tag{10}$$

If we next substitute Eq. (6) with Eqs. (7), (9) and (10) into Eq. (5) and balance z''' with  $z^n z'$ , we have

$$M+3 = nM + M + 1 \text{ resulting in } M = \frac{2}{n}.$$
 (11)

In order to obtain the positive integer M values for Eq. (6), we use the transformation

$$u = v^{\frac{z}{n}}$$
(12)

in Eq. (5), which yields

$$2wn^{2}v^{2}v' + 2\varepsilon kn^{2}v^{4}v' + 2\mu k^{3}n^{2}v^{2}v''' - 6\mu k^{3}n^{2}vv'v'' + 12\mu k^{3}nvv'v'' + 4\mu k^{3}n^{2}(v')^{3} - 12\mu k^{3}n(v')^{3} -rn^{3}v^{3} + rn^{3}v^{5} = 0.$$
(13)

If we now apply the same procedure by using the expression for v as

$$v = \sum_{i=0}^{M} a_i z^i \tag{14}$$

in Eq. (13) with Eqs. (7), (9) and (10), and balancing the highest power of z, we have M = 1, which yields the solution of Eq. (13) to be in the form

$$v = a_0 + a_1 z \,. \tag{15}$$

Let us next use Eq. (15) in Eq. (13) and collect the coefficients for the same powers of z as follows:

$$z^{0}: 4a_{1}^{3}k^{3}\mu n^{2} - 4k^{3}\mu n^{2}a_{0}^{2}a_{1} + 2kn^{2}\varepsilon a_{0}^{4}a_{1} + n^{3}ra_{0}^{5}$$
  
$$-12a_{1}^{3}k^{3}\mu n + 8a_{1}^{3}k^{3}\mu - n^{3}ra_{0}^{3} + 2n^{2}wa_{0}^{2}a_{1} = 0,$$
  
$$z^{1}: 4k^{3}\mu n^{2}a_{0}a_{1}^{2} + 8kn^{2}\varepsilon a_{0}^{3}a_{1}^{2} + 5n^{3}ra_{0}^{4}a_{1} + 4n^{2}wa_{0}a_{1}^{2}$$
  
$$-24k^{3}\mu na_{0}a_{1}^{2} - 3n^{3}ra_{0}^{2}a_{1} = 0,$$

$$z^{2}: 16k^{3}\mu n^{2}a_{0}^{2}a_{1} - 4a_{1}^{3}k^{3}\mu n^{2} - 2kn^{2}\varepsilon a_{0}^{4}a_{1} - 24a_{1}^{3}k^{3}\mu + 12kn^{2}\varepsilon a_{0}^{2}a_{1}^{3} + 10n^{3}ra_{0}^{3}a_{1}^{2} + 12a_{1}^{3}k^{3}\mu n + 2n^{2}wa_{1}^{3} - 3n^{3}ra_{0}a_{1}^{2} - 2n^{2}wa_{0}^{2}a_{1} = 0,$$
(16)

$$z^{3}: 8k^{3}\mu n^{2}a_{0}a_{1}^{2} - 8kn^{2}\varepsilon a_{0}^{3}a_{1}^{2} + 8kn^{2}\varepsilon a_{0}a_{1}^{4} - n^{3}ra_{1}^{3}$$
  
+10n<sup>3</sup>ra\_{0}^{2}a\_{1}^{3} + 48k^{3}\mu na\_{0}a\_{1}^{2} - 4n^{2}wa\_{0}a\_{1}^{2} = 0,

$$\begin{split} z^4: & 4a_1^3k^3\mu n^2 - 12k^3\mu n^2a_0^2a_1 - 12kn^2\varepsilon a_0^4a_1^3 - 2n^2wa_1^3 \\ & + 2kn^2\varepsilon a_1^5 + 5n^3ra_0a_1^4 + 12a_1^3k^3\mu n + 24a_1^3k^3\mu = 0 \,, \\ z^5: & n^3ra_1^5 - 12k^3\mu n^2a_0a_1^2 - 8kn^2\varepsilon a_0a_1^4 - 24k^3\mu na_0a_1^2 \,, \\ z^6: & -4a_1^3k^3\mu n^2 - 2kn^2\varepsilon a_1^5 - 12a_1^3k^3\mu n - 8a_1^3k^3\mu = 0 \,. \end{split}$$

As the final step, solving the nonlinear system (16) for  $a_0$ ,  $a_1$  and nonzero k, w,  $\varepsilon$ ,  $\mu$ , r, n parameters by using Maple, we have the following solutions for  $\varepsilon\mu < 0$ :

$$u_{1,2} = \left(\frac{\tanh(kx + wt + \xi_0) \pm 1}{2}\right)^{\frac{2}{n}}$$
(17)

and

$$u_{3,4} = \left(\frac{\coth\left(kx + wt + \xi_0\right) \pm 1}{2}\right)^{\frac{2}{n}},\tag{18}$$

where

$$r = \mp \frac{\varepsilon (n+4)\sqrt{-2\varepsilon \mu (n+1)(n+2)}}{2\mu (n+1)^2 (n+2)},$$

$$k = \pm \frac{r(n+1)n}{2\varepsilon (n+4)},$$

$$w = \pm \frac{r(n^2 + 6n + 12)n}{4(n+2)(n+4)}.$$
(19)

Figure 1 exhibits the long-time behaviour for the global solution  $u_1(x,t)$  of the gKdV-Fisher equation in Eq. (17) for the different values of the parameters, which represent kink and antikink waves. One can easily see that the propagations of waves are backward in Figure 1-(b), i.e. k, w > 0, and forward in Figure 1-(a) and Figure 1-(c). Another global solution  $u_2(x,t)$  in Eq. (17) is displayed in Figure 2 for the given parameters, which also exhibit kink and antikink waves.

On the other hand, the blow-up solutions  $u_3(x,t)$  and  $u_4(x,t)$  in Eq. (18) are displayed in Figure 3 and Figure 4, respectively.



(a)  $\varepsilon = -1$ ,  $\mu = 1$ ,  $r = 5\sqrt{3}/12$ , n = 1,  $\xi_0 = 0$ .



**(b)**  $\varepsilon = 1$ ,  $\mu = -1$ ,  $r = 5\sqrt{3}/12$ , n = 1,  $\xi_0 = 0$ .



**Figure 1.** The solution  $u_1(x,t)$  of the gKdV-Fisher equation for the different suitable parameters.

Because of the nature of the complex nonlinear phenomena, it is reasonable to find the blow-up solutions for the mix of the different entities with nonlinearity. Fortunately, we can reveal the cooperative combinations of dispersion, convection and reaction with the parameters in Eq. (19) for the global solutions of the gKdV-Fisher equation, which represent kink and antikink waves, see Figure 1 and 2.



(a) 
$$\varepsilon = -1$$
,  $\mu = 1$ ,  $r = -5\sqrt{3}/12$ ,  $n = 1$ ,  $\xi_0 = 0$ .



(b)  $\varepsilon = -1$ ,  $\mu = 1$ ,  $r = -\sqrt{6}/6$ , n = 2,  $\xi_0 = 0$ . Figure 2. The solution  $u_2(x,t)$  of the gKdV-Fisher equation

for the different suitable parameters.







**(b)**  $\varepsilon = -1$ ,  $\mu = 1$ ,  $r = \sqrt{6}/6$ , n = 2,  $\xi_0 = 0$ .

**Figure 3.** The solution  $u_3(x,t)$  of the gKdV-Fisher equation for the different suitable parameters.



(a)  $\varepsilon = -1$ ,  $\mu = 1$ ,  $r = -5\sqrt{3}/12$ , n = 1,  $\xi_0 = 0$ .



**(b)**  $\varepsilon = -1$ ,  $\mu = 1$ ,  $r = -\sqrt{6}/6$ , n = 2,  $\xi_0 = 0$ .

**Figure 4.** The solution  $u_4(x,t)$  of the gKdV-Fisher equation for the different suitable parameters.

## 3. Conclusion

A nonlinear dispersion-convection-reaction model, called the gKdV-Fisher equation, has been introduced to investigate the travelling wave solutions. A classical and efficient the Riccati equation method has been used to investigate two new global and two new blow-up solutions. One can easily see that the reaction term in the proposed equation yields kink and antikink wave solutions, which can be used in the other various numerical and analytical investigations on the application of such combined model to scientific problems. Further research would be based on investigating N-soliton solutions of the third and higher odd-order PDEs including Fisher-type nonlinearity.

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